# Conditional expectation and Bayes' rule for quantum random variables and positive operator valued measures

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References

(F-P-S) Classical and nonclassical randomness in quantum measurements by Douglas Farenick, Sarah Plosker, and Jerrod Smith. *J. Math. Phys.*, 52:122204, 2011.

**(F-K)** Conditional expectation and Bayes rule for quantum random variables and positive operator valued measures by Douglas Farenick and MJK. J. Math. Phys., 53:042201, 2012.

Background

A measurement of a quantum system is represented mathematically by a positive operator valued measure  $\nu$  which is defined on a  $\sigma$ -algebra  $\mathcal{O}(X)$  of measurement events such that whenever a measurement is made with the system in state  $\rho$ , the measurement event  $E \in \mathcal{O}(X)$  will occur with probability

$$\operatorname{Tr}(\rho\nu(E)).$$

Reference. The Quantum Theory of Measurement by Busch, Lahti, Mittelstaedt, LNP, Springer, 1991.

In practice, quantum measurements of an actual physical system are made by way of some apparatus and so X is often assumed to by finite.

Mathematically, however, there is no need for such a restriction and so one of our goals is to approach the theory of quantum measurement under the assumption that X be arbitrary.

#### Some notation

X, a locally compact Hausdorff space

 $\mathcal{O}(X)$ , the Borel  $\sigma$ -algebra of subsets of X

 $\mathcal{F}(X)$ , a sub- $\sigma$ -algebra of  $\mathcal{O}(X)$ 

 $\mathcal{H}$ , a d-dimensional Hilbert space

 $\mathcal{B}(\mathcal{H})$ , the space of (bounded) linear operators on  $\mathcal{H}$ 

 $\mathrm{Tr}:\mathcal{B}(\mathcal{H}) \to \mathbb{C}$ , the canonical trace functional

 $\mathcal{B}(\mathcal{H})_+ = \{a \in \mathcal{B}(\mathcal{H}) : \langle a\zeta, \zeta \rangle \geq 0 \ \forall \ \zeta \in \mathcal{H} \}$ , the space of positive operators

 $S(\mathcal{H})$ , the state space of  $\mathcal{H}$ , namely the set of all density operators  $\rho \in \mathcal{B}(\mathcal{H})_+$  with  $Tr(\rho)=1$ 

 $\mathrm{Eff}(\mathcal{H})$ , the set of quantum effects, namely those positive operators  $h \in \mathcal{B}(\mathcal{H})_+$  such that every eigenvalue of h lies in [0,1].

$$S(\mathcal{H}) \subset Eff(\mathcal{H})$$

#### Positive operator valued measures

A set function  $\nu: \mathcal{F}(X) \to \mathcal{B}(\mathcal{H})$  is called a positive operator valued measure on  $(X,\mathcal{F}(X))$  if

- 1.  $\nu(E) \in \text{Eff}(\mathcal{H})$  for every  $E \in \mathcal{F}(X)$ ,
- 2.  $\nu(X) \neq 0$ , and
- 3. for every countable collection  $\{E_k\}_{k\in\mathbb{N}}\subseteq\mathcal{F}(X)$  with  $E_j\cap E_k=\emptyset$  for  $j\neq k$  we have

$$\nu\left(\bigcup_{k\in\mathbb{N}} E_k\right) = \sum_{k\in\mathbb{N}} \nu(E_k)$$

where the convergence on the right side of the previous equality is with respect to the  $\sigma$ -weak topology of  $\mathcal{B}(\mathcal{H})$ .

If  $\nu(X) = 1 \in \mathcal{B}(\mathcal{H})$ , we call it a positive operator valued probability measure.

**Notation.**  $POVM_{\mathcal{H}}(X)$  or  $POVM_{\mathcal{H}}^1(X)$ 

#### Quantum random variables

A quantum random variable on X is a function  $\psi:X\to\mathcal{B}(\mathcal{H})$  such that

$$x \mapsto \operatorname{Tr}(\rho \psi(x))$$

is a complex random variable on X for every density operator  $\rho \in S(\mathcal{H})$ .

**Recall.** A random variable is a Borel-measurable function on a measure space X with  $\mu(X)=1$ .

Quantum averaging (F-P-S, F-K)

**Theorem.** There is a definition of integral whereby a quantum random variable  $\psi$  may be integrated against the positive operator valued probability measure  $\nu$  to produce an operator

$$\mathbb{E}_{\nu} [\psi] := \int_{X} \psi \, d\nu \in \mathcal{B}(\mathcal{H}).$$

**Example.** If  $X = \{x_1, x_2, \dots, x_n\}$ , then

$$\mathbb{E}_{\nu} \left[ \psi \right] = \sum_{j=1}^{n} h_{j}^{1/2} \psi(x_{j}) h_{j}^{1/2}$$

where  $h_j = \nu(x_j)$ .

#### The principal Radon-Nikodým derivative (F-P-S)

Let  $\nu \in \text{POVM}_{\mathcal{H}}(X, \mathcal{F}(X))$  so that  $\mu(E) = \frac{\text{Tr}(\nu(E))}{d}$  is a Borel measure.

Let  $\{e_1, \ldots, e_d\}$  be an orthonormal basis of  $\mathcal{H}$ .

Let  $\nu_{ij}: \mathcal{F}(X) \to \mathbb{C}$  be defined by  $\nu_{ij}(E) = \langle \nu(E)e_j, e_i \rangle$  so that  $\nu_{ij} \ll_{\mathrm{ac}} \mu$ . By classical R-N Theorem, there exists a unique function

$$\frac{\mathrm{d}\nu_{ij}}{\mathrm{d}\mu} \in L^1(X, \mathcal{F}(X), \mu)$$

such that

$$\nu_{ij}(E) = \int_E \frac{\mathrm{d}\nu_{ij}}{\mathrm{d}\mu} \,\mathrm{d}\mu.$$

The function

$$\frac{\mathrm{d}\nu}{\mathrm{d}\mu} = \sum_{i,j=1}^{d} \frac{\mathrm{d}\nu_{ij}}{\mathrm{d}\mu} \otimes e_{ij}$$

where  $e_{ij} \in \mathcal{B}(\mathcal{H})$  sends  $e_j$  to  $e_i$  and  $e_k$  to 0 is called the principal Radon-Nikodým derivative of  $\nu$ .

## The non-principal Radon-Nikodým derivative (F-P-S)

**Theorem.** If  $\nu_1, \nu_2 \in \text{POVM}_{\mathcal{H}}(X, \mathcal{F}(X))$ , then the following statements are equivalent.

- 1.  $\nu_2 \ll_{ac} \nu_1$ , i.e.,  $\nu_2(E) = 0$  whenever  $\nu_1(E) = 0$ .
- 2. There exists a bounded  $\nu_1$ -integrable  $\mathcal{F}(X)$ -measurable function  $\varphi:(X,\mathcal{F}(X))\to\mathcal{B}(\mathcal{H})$ , unique up to sets of  $\nu_1$ -measure zero, such that

$$\nu_2(E) = \int_E \varphi \, \mathrm{d}\nu_1 \tag{1}$$

for every  $E \in \mathcal{F}(X)$ .

Moreover, if the equivalent conditions above hold and if  $\mu_j = \mu_{\nu_j}$  is the finite Borel measure induced by  $\nu_j$ , then  $\mu_2 \ll_{\rm ac} \mu_1$  and

$$\varphi = \left(\frac{\mathrm{d}\mu_2}{\mathrm{d}\mu_1}\right) \left[ \left(\frac{\mathrm{d}\nu_1}{\mathrm{d}\mu_1}\right)^{-1/2} \left(\frac{\mathrm{d}\nu_2}{\mathrm{d}\mu_2}\right) \left(\frac{\mathrm{d}\nu_1}{\mathrm{d}\mu_1}\right)^{-1/2} \right]$$
(2)

i.e.,  $\varphi$  is the non-principal Radon-Nikodým derivative of  $\nu_2$  wrt  $\nu_1$  so we write

$$\frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} = \varphi.$$

### A non-commutative multiplication (F-K)

If  $a, b \in \mathcal{B}(\mathcal{H})_+$  are both invertible, then the geometric mean of a and b is defined by setting

$$a\#b = a^{1/2}(a^{-1/2}ba^{-1/2})^{1/2}a^{1/2}.$$

**Definition.** Suppose that  $\nu_1$ ,  $\nu_2 \in \text{POVM}^1_{\mathcal{H}}(X)$  with  $\nu_2 \ll_{\text{ac}} \nu_1$ , and let  $\mu_j = \mu_{\nu_j}$  be the induced Borel probability measures. If  $\psi: X \to \mathcal{B}(\mathcal{H})$  is a quantum random variable, then

$$\psi \boxtimes \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} = \left( \left( \frac{\mathrm{d}\nu_1}{\mathrm{d}\mu_1} \right)^{-1} \# \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} \right) \left( \frac{\mathrm{d}\nu_1}{\mathrm{d}\mu_1} \right)^{1/2} \psi \left( \frac{\mathrm{d}\nu_1}{\mathrm{d}\mu_1} \right)^{1/2} \left( \left( \frac{\mathrm{d}\nu_1}{\mathrm{d}\mu_1} \right)^{-1} \# \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} \right).$$

**Remark.** In the commutative setting—and, in particular, in the classical case of  $\mathcal{H}=\mathbb{C}$ —the multiplication defined by  $\boxtimes$  reduces to ordinary multiplication. That is, if  $a,b\in\mathcal{B}(\mathcal{H})_+$  commute, then  $a\#b=a^{1/2}b^{1/2}=b^{1/2}a^{1/2}=b\#a$ . Thus, if  $\psi$ ,  $\frac{\mathrm{d}\nu_1}{\mathrm{d}\mu_1}$ , and  $\frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1}$  are pairwise commuting, then

$$\psi \boxtimes \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} = \psi \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} = \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} \psi.$$

Change of quantum measurement (F-K)

**Theorem.** Suppose that  $\nu_1$ ,  $\nu_2 \in \mathrm{POVM}^1_{\mathcal{H}}(X)$  with  $\nu_2 \ll_{\mathrm{ac}} \nu_1$ , and let  $\mu_j = \mu_{\nu_j}$  be the induced Borel probability measures. If  $\psi: X \to \mathcal{B}(\mathcal{H})$  is a  $\nu_2$ -integrable quantum random variable, then

$$\psi \boxtimes \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1}$$

is a  $u_1$ -integrable quantum random variable and

$$\mathbb{E}_{\nu_2} \left[ \psi \right] = \mathbb{E}_{\nu_1} \left[ \psi \boxtimes \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} \right]$$

or, equivalently,

$$\int_X \psi \, \mathrm{d}\nu_2 = \int_X \psi \boxtimes \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} \, \mathrm{d}\nu_1.$$

Theorems for the Radon-Nikodým derivatives (F-K)

Theorem (Chain Rule). If  $\nu_1, \nu_2, \nu_3 \in \text{POVM}^1_{\mathcal{H}}(X)$  with  $\nu_1 \ll_{\text{ac}} \nu_2 \ll_{\text{ac}} \nu_3$ , then

$$\frac{\mathrm{d}\nu_1}{\mathrm{d}\nu_2} \boxtimes \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_3} = \frac{\mathrm{d}\nu_1}{\mathrm{d}\nu_3}.$$

Corollary. If  $\nu_1, \nu_2 \in \text{POVM}^1_{\mathcal{H}}(X)$  with  $\nu_2 \ll_{\text{ac}} \nu_1$  and  $\nu_1 \ll_{\text{ac}} \nu_2$ , then

$$\frac{\mathrm{d}\nu_1}{\mathrm{d}\nu_2} \boxtimes \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} = \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} \boxtimes \frac{\mathrm{d}\nu_1}{\mathrm{d}\nu_2} = 1.$$

Quantum conditional expectation (F-K)

**Theorem.** Suppose that  $\nu \in \mathrm{POVM}^1_{\mathcal{H}}(X)$  and that  $\psi : X \to \mathcal{B}(\mathcal{H})_+$  is a  $\nu$ -integrable quantum random variable with  $\mathbb{E}_{\nu} \left[ \psi \right] \neq 0$ . If  $\mathcal{F}(X)$  is a sub- $\sigma$ -algebra of  $\mathcal{O}(X)$ , then there exists a function  $\varphi : X \to \mathcal{B}(\mathcal{H})$  such that

- 1.  $\varphi$  is  $\mathcal{F}(X)$ -measurable,
- 2.  $\varphi$  is  $\nu$ -integrable, and
- 3.  $\mathbb{E}_{\nu} [\psi \chi_E] = \mathbb{E}_{\nu} [\varphi \chi_E]$  for every  $E \in \mathcal{F}(X)$ .

Moreover, if  $\tilde{\varphi}$  is any other  $\nu$ -integrable  $\mathcal{F}(X)$ -measurable function satisfying  $\mathbb{E}_{\nu}\left[\psi\chi_{E}\right]=\mathbb{E}_{\nu}\left[\tilde{\varphi}\chi_{E}\right]$  for every  $E\in\mathcal{F}(X)$ , then  $\nu(\{x\in X:\varphi(x)\neq\tilde{\varphi}(x)\})=0$ .

We write

$$\varphi = \mathbb{E}_{\nu} \left[ \psi | \mathcal{F}(X) \right].$$

Quantum Bayes' rule (F-K)

**Theorem.** Let  $\nu_1$ ,  $\nu_2 \in \mathrm{POVM}^1_{\mathcal{H}}(X)$  with  $\nu_2 \ll_{\mathrm{ac}} \nu_1$  and  $\nu_1 \ll_{\mathrm{ac}} \nu_2$ .

If  $\psi: X \to \mathcal{B}(\mathcal{H})_+$  is a quantum random variable with  $\mathbb{E}_{\nu_2} [\psi] \neq 0$  and  $\mathcal{F}(X)$  is a sub- $\sigma$ -algebra of  $\mathcal{O}(X)$ , then

$$\mathbb{E}_{\nu_2} \left[ \psi | \mathcal{F}(X) \right] \boxtimes \mathbb{E}_{\nu_1} \left[ \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} \middle| \mathcal{F}(X) \right] = \mathbb{E}_{\nu_1} \left[ \psi \boxtimes \frac{\mathrm{d}\nu_2}{\mathrm{d}\nu_1} \middle| \mathcal{F}(X) \right].$$